

# Universal quantum groups acting on classical and quantum spaces

## Lecture 1 - Compact quantum groups and their actions

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# Classical symmetry groups

Groups entered mathematics as collections of symmetries of a given object (a finite set, a figure on the plane, a manifold, a space of solutions of an equation).

How do we define a symmetry group of a given object  $X$ ? We look at the family of 'all possible transformations' of  $X$ , usually preserving some structure of  $X$ .

In modern language: we search for a *universal* object in the category of all groups acting on  $X$ .

# Compact quantum groups

We will be interested in the cases when the classical symmetry groups are compact/finite.

The idea of defining a compact quantum group is based on 'quantising' the algebra of continuous functions on a compact group.

## Definition ([Wor<sub>2</sub>])

A **compact quantum group** is a unital  $C^*$ -algebra  $A$  with a unital  $*$ -algebra homomorphism  $\Delta : A \rightarrow A \otimes A$  such that

$$(\text{id} \otimes \Delta) \circ \Delta = (\Delta \otimes \text{id}) \circ \Delta \quad (\text{coassociativity})$$

and

$$\overline{\Delta(A)(A \otimes 1_A)} = A \otimes A = \overline{\Delta(A)(1_A \otimes A)} \quad (\text{quantum cancellation rules}).$$

The tensor products here are minimal/spatial tensor products of  $C^*$ -algebras.

Let us list some basic properties of and notions related to compact quantum groups we will need:

- there exists a unique bi-invariant state, a so called *Haar state*  $h \in A^*$ :

$$(h \otimes \text{id}_A) \circ \Delta = (\text{id}_A \otimes h) \circ \Delta = h(\cdot)1_A.$$

- a *finite-dimensional unitary corepresentation* of  $A$  is a unitary matrix  $U \in M_n \otimes A$  such that (in the leg notation)

$$(\text{id}_{M_n} \otimes \Delta)(U) = U_{12}U_{13}.$$

In terms of the entries:

$$\Delta(U_{ij}) = \sum_{k=1}^n U_{ik} \otimes U_{kj};$$

- $A$  admits a unique dense Hopf\*-algebra  $(\mathcal{A}, \Delta, \epsilon, S)$ , spanned by the *coefficients* of the finite-dimensional unitary corepresentations of  $A$ ; in particular  $\Delta : \mathcal{A} \rightarrow \mathcal{A} \odot \mathcal{A}$ . Hopf\*-algebras arising in this way admit an intrinsic characterisation, which makes the theory of compact quantum groups in a sense completely algebraic.

- $h$  is faithful on  $\mathcal{A}$ ; if it is faithful on  $A$ , we say that  $A$  is in the reduced form. This can be always achieved by considering the  $C^*$ -completion of the GNS-representation  $\pi_h$  of  $\mathcal{A}$ . Conversely, one can consider the universal  $C^*$ -completion of  $\mathcal{A}$  - if the reduced and universal object coincide, we say that  $A$  is *coamenable*.
- the coproduct has a unique normal extension to a coassociative unital  $*$ -homomorphism  $\tilde{\Delta} : M \rightarrow M \overline{\otimes} M$ , where  $M = \pi_h(A)''$ ;
- if  $h$  is tracial, then  $(A, \Delta)$  is a *Kac algebra of compact type*; in particular  $S^2 = \text{id}_{\mathcal{A}}$ .

# Examples of compact quantum groups

- commutative CQGs - algebras  $C(G)$ , where  $G$ -compact group. The coproduct is determined by (recall that  $C(G) \otimes C(G) \approx C(G \times G)$ )

$$\Delta(f)(s, t) = f(st), \quad s, t \in G, f \in C(G).$$

- cocommutative CQGs - essentially (subtlety is related to a possible non-coamenability) algebras  $C^*(\Gamma)$  or  $C_r^*(\Gamma)$ , where  $\Gamma$  - discrete group. The coproduct is given by the (continuous linear extension of)

$$\Delta(\pi_\gamma) = \pi_\gamma \otimes \pi_\gamma, \quad \gamma \in \Gamma.$$

## Examples continued: deformations of classical compact Lie groups

- Note that  $C(SU(2))$  is a commutative unital  $C^*$ -algebra generated by the functions  $\alpha, \gamma : SU(2) \rightarrow \mathbb{C}$  such that

$$\alpha^* \alpha + \gamma^* \gamma = 1.$$

Group multiplication on  $SU(2)$  induces on  $C(SU(2))$  the coproduct

$$\Delta(\alpha) = \alpha \otimes \alpha - \gamma^* \otimes \gamma, \quad \Delta(\gamma) = \gamma \otimes \alpha + \alpha^* \otimes \gamma.$$

Let  $q \in [-1, 1] \setminus \{0\}$ . Define  $SU_q(2)$  - unital  $C^*$ -algebra generated by operators  $\alpha, \gamma$  such that:

$$\begin{aligned} \alpha^* \alpha + \gamma^* \gamma &= 1, & \alpha \alpha^* + q^2 \gamma^* \gamma &= 1, \\ \gamma^* \gamma &= \gamma \gamma^*, & q \gamma \alpha &= \alpha \gamma, & q \gamma^* \alpha &= \alpha \gamma^*. \end{aligned}$$

The coproduct making  $SU_q(2)$  a compact quantum group is given by the formulas

$$\Delta(\alpha) = \alpha \otimes \alpha - q \gamma^* \otimes \gamma, \quad \Delta(\gamma) = \gamma \otimes \alpha + \alpha^* \otimes \gamma.$$

# Compact matrix quantum groups ([Wor<sub>1</sub>])

If a compact quantum group  $A$  admits  $n \in \mathbb{N}$  and a unitary matrix  $U = (u_{ij})_{i,j=1}^n \in M_n(A)$  such that

- $\Delta(u_{ij}) = \sum_{k=1}^n u_{ik} \otimes u_{kj}$ ;
- $\{u_{ij} : i, j = 1, \dots, n\}$  generates  $A$  as a  $C^*$ -algebra,

then  $A$  is called a *compact matrix quantum group* and  $U$  a *fundamental unitary corepresentation* of  $A$ . The counit of  $A$  is determined by the formula  $\epsilon(u_{ij}) = \delta_{ij}$ , the antipode by the formula  $S(u_{ij}) = u_{ji}^*$ .

# Actions

The action of a group  $G$  on a set  $X$  can be described as a map  $\alpha : G \times X \rightarrow X$  satisfying certain natural conditions. In the quantum case we (as usual) ‘invert the arrows’:

## Definition ([Pod], [Wan<sub>2</sub>])

Let  $A$  be a compact quantum group and let  $B$  be a unital  $C^*$ -algebra. A map

$$\alpha : B \rightarrow A \otimes B$$

is called a (left, continuous) action of  $A$  on  $B$  if  $\alpha$  is a unital  $*$ -homomorphism,

$$(\Delta \otimes \text{id}_B) \circ \alpha = (\text{id}_A \otimes \alpha) \circ \alpha$$

and additionally  $\alpha(B)(A \otimes 1_B)$  is dense in  $A \otimes B$  (*Podleś/continuity condition*).

Unless stated otherwise when we say that  $\alpha : B \rightarrow A \otimes B$  is an action we mean that all the above are satisfied. Note that each compact quantum group acts on itself via the coproduct.

## On *continuity* of actions and its consequences

If  $A = C(G)$  for a compact group  $G$ ,  $X$  is a compact space and  $G \curvearrowright X$  is a continuous action, then in the quantum picture we define

$$\alpha : C(X) \rightarrow C(G) \otimes C(X) \approx C(G \times X)$$

via

$$\alpha(f)(g, x) = f(gx), \quad g \in G, x \in X, f \in C(X).$$

Conversely given an action of  $\alpha : C(X) \rightarrow C(X) \otimes C(G)$  we define for each  $g \in G$  first

$$\alpha_g = (\text{ev}_g \otimes \text{id}_{C(X)}) \circ \alpha$$

and then note that the 'action relation' implies that  $\alpha_g \circ \alpha_h = \alpha_{gh}$  for  $g, h \in G$ .

The condition  $\alpha(B)(A \otimes 1_B)$  is dense in  $A \otimes B$  is a kind of a nondegeneracy or continuity property for the action, in the case above guaranteeing that each  $\alpha_g$  is a surjection. It excludes for example  $\alpha = \rho(\cdot)1_A \otimes 1_B$ , where  $\rho$  is a character on  $B$ .

The continuity condition implies that there exists  $\mathcal{B}$ , a dense  $*$ -subalgebra of  $B$ , such that  $\alpha|_{\mathcal{B}} : \mathcal{B} \rightarrow \mathcal{A} \odot \mathcal{B}$  is a Hopf $*$ -algebraic action ([Wan<sub>2</sub>], [Boc]); in particular we have

$$(\epsilon \otimes \text{id}_{\mathcal{B}})\alpha|_{\mathcal{B}} = \text{id}_{\mathcal{B}}.$$

The algebra  $\mathcal{B}$  is a linear span of *spectral subspaces*  $\mathcal{B}_u$  corresponding to *irreducible* unitary corepresentations of  $A$ .

The action  $\alpha : B \rightarrow A \otimes B$  is called *faithful* if the set  $\text{Lin}\{(\text{id}_A \otimes \phi) \circ \alpha(b) : b \in B, \phi \in B^*\}$  is dense in  $A$ .

Exercises:

- Check that if  $U \in M_n(A) \approx A \otimes M_n$  is a unitary corepresentation, then the map

$$x \rightarrow U^*(1_A \otimes x)U, \quad x \in M_n,$$

defines an action of  $A$  on  $M_n$ .

- Arbitrary (not necessarily finite-dimensional) unitary corepresentations are defined as unitaries in the multiplier algebra  $M(A \otimes K(H))$ , where  $H$  is a Hilbert space, again satisfying the suitably understood relation

$$(\Delta \otimes \text{id}_{K(H)})(U) = U_{13}U_{23};$$

do they define actions of  $A$  in our sense by the formula from the first part of the exercise?

# Lifting of actions to von Neumann algebras

## Definition

We say that the action  $\alpha : B \rightarrow A \otimes B$  preserves a functional  $\omega \in B^*$  if

$$\forall_{b \in B} (\text{id}_A \otimes \omega)(\alpha(b)) = \omega(b)1_A.$$

## Theorem ([Wan<sub>2</sub>])

Let  $\alpha : B \rightarrow A \otimes B$  be an action preserving a state  $\tau \in B^*$ . Then  $\alpha$  lifts to a normal  $*$ -homomorphism  $\tilde{\alpha} : \pi_\tau(B)'' \rightarrow \pi_h(A)'' \overline{\otimes} \pi_\tau(B)''$  defined by the formula

$$\tilde{\alpha}(\pi_\tau(b)) = (\pi_h \otimes \pi_\tau)(\alpha(b)), \quad b \in B,$$

which is an action in the Hopf-von Neumann algebra sense.

# Ergodic actions

## Definition

We say that the action  $\alpha : B \rightarrow A \otimes B$  is ergodic if the fixed point space  $\text{Fix}(\alpha) = \{b \in B : \alpha(b) = 1_A \otimes b\}$  is one-dimensional.

For arbitrary action  $\alpha$  the map  $E_\alpha : B \rightarrow B$  defined by

$$E_\alpha = (h \otimes \text{id}_B)\alpha$$

is a norm-one projection from  $B$  onto  $\text{Fix}(\alpha)$ . Moreover  $\{b \in B : \alpha(b) = 1_A \otimes b\}$  is dense in  $\text{Fix}(\alpha)$ .

## Theorem ([Boc], [BDRV])

*Spectral subspaces of ergodic actions are finite-dimensional, with their dimensions dominated by the (quantum) dimensions of the corresponding corepresentations.*

Some ergodic actions can be described in abstract categorical terms. The theory becomes much richer than the classical one.

## Category of CQGs acting on a given $C^*$ -algebra

Consider the category  $\mathfrak{C}(B) := \{(A, \alpha)\}$  of compact quantum groups acting on a given  $C^*$ -algebra  $B$ . A morphism in the category  $\mathfrak{C}(B)$ :

$$\gamma : (A_1, \alpha_1) \rightarrow (A_2, \alpha_2)$$

is a unital  $*$ -homomorphism  $\gamma : A_1 \rightarrow A_2$  such that

$$(\gamma \otimes \gamma) \circ \Delta_{A_1} = \Delta_{A_2} \circ \gamma, \quad \alpha_2 = (\gamma \otimes \text{id}_B) \circ \alpha_1.$$

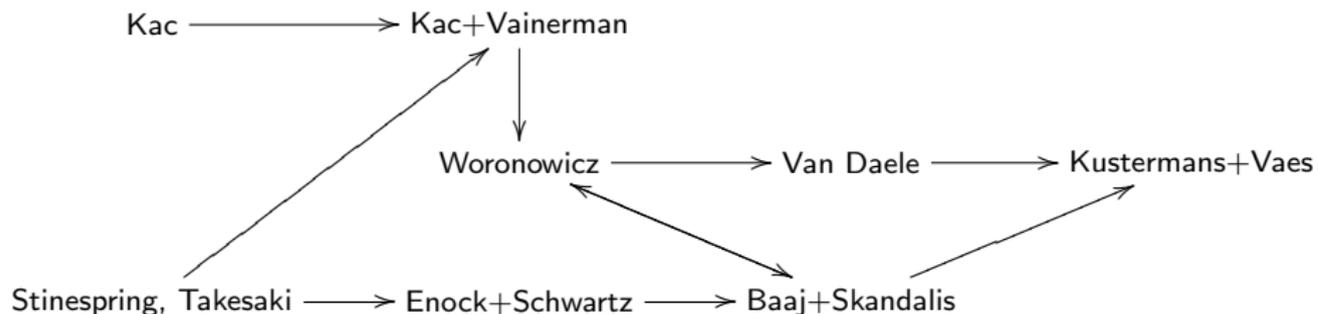
We say that the category  $\mathfrak{C}(B)$  admits a universal object, if there is  $(A_u, \alpha_u)$  in  $\mathfrak{C}(B)$  such that for all  $(A, \alpha)$  in  $\mathfrak{C}(B)$  there exists a unique morphism  $\gamma : (A_u, \alpha_u) \rightarrow (A, \alpha)$ . Abstract categorical nonsense implies that if such a universal object exists, it is unique.

$(A_u, \alpha_u)$  – quantum symmetry group of  $B$

Further we will also consider categories of actions preserving some ‘extra’ structure of  $B$ .

# Locally compact quantum semigroups

The story of development of locally compact quantum groups is long and complicated - we will just offer a diagram:



# Locally compact quantum semigroups

For now it suffices for us to work in a simplified context of locally compact quantum semigroups:

## Definition

A **locally compact quantum semigroup** is a  $C^*$ -algebra  $A$  equipped with a nondegenerate  $*$ -homomorphism  $\Delta : A \rightarrow M(A \otimes A)$  which is coassociative:

$$(\text{id}_A \otimes \Delta) \circ \Delta = (\Delta \otimes \text{id}_A) \circ \Delta.$$

Some comments on the definition:

- The multiplier algebra of a  $C^*$ -algebra  $C$  is the largest  $C^*$ -algebra containing  $C$  as an essential ideal. It is equipped with the *strict* topology: locally convex topology determined by the seminorms  $\{l_c, r_c : c \in C\}$ , where  $l_c(x) = \|cx\|$ ,  $r_c(x) = \|xc\|$  ( $x \in M(C)$ ).
- A bounded linear map  $\phi : C \rightarrow M(D)$  is called *strict* ([Kus]) if it is strictly continuous on bounded subsets of  $C$ . Such a map admits a unique strict extension to  $\tilde{\phi} : M(C) \rightarrow M(D)$ .
- A  $*$ -homomorphism  $\pi : C \rightarrow M(D)$  is called nondegenerate if  $\pi(C)D$  is dense in  $D$ . Nondegenerate  $*$ -homomorphisms are strict.
- When we compose strict maps we always have in mind respective strict extensions: so for example it makes sense to write  $(\text{id}_A \otimes \Delta) \circ \Delta$  and understand it as a map from  $A$  (or  $M(A)$ ) to  $M(A \otimes A \otimes A)$ .

# 'Non-compact' actions

## Definition

Let  $(A, \Delta)$  be a locally compact quantum semigroup and let  $B$  be a  $C^*$ -algebra. An action of  $A$  on  $B$  is a nondegenerate  $*$ -homomorphism  $\alpha : B \rightarrow M(A \otimes B)$  such that

$$(\Delta \otimes \text{id}_B) \circ \alpha = (\text{id}_A \otimes \alpha) \circ \alpha,$$

$\text{Lin}\{(a \otimes 1_{M(B)})\alpha(b) : a \in A, b \in B\}$  is (contained in and) dense in  $A \otimes B$ .

Exercises ([Sof<sub>2</sub>]):

- Let  $G$  be a locally compact group and  $B$  a  $C^*$ -algebra. Put  $A := C_0(G)$ . Assume  $\alpha : B \rightarrow M(A \otimes B)$  is a nondegenerate  $*$ -homomorphism such that  $(\Delta \otimes \text{id}_B) \circ \alpha = (\text{id}_A \otimes \alpha) \circ \alpha$ . Prove that the following are equivalent:
  - $(\text{ev}_e \otimes \text{id}_B) \circ \alpha = \text{id}_B$ ;
  - $\text{Ker}(\alpha) = \{0\}$ ;
  - $\text{Lin}\{(a \otimes 1_{M(B)})\alpha(b) : a \in A, b \in B\}$  is dense in  $A \otimes B$ ;
  - $\text{Lin}\{(\omega_f \otimes \text{id}_B)\alpha(b) : f \in L^1(G), b \in B\}$  is dense in  $B$  ( $\omega_f \in C_0(G)^*$  is given by the integration with respect to the Haar measure on  $G$ ).
- Show that if the above conditions hold,  $\alpha$  induces a continuous action of  $G$  on  $B$  by automorphisms.

Even when we pass to locally compact quantum *groups* (i.e. assume the existence of suitably well-behaved left- and right- invariant weights), there is still no ‘nice’ underlying algebraic object  $\mathcal{A}$  behind our  $C^*$ -algebra (as it was in the compact case). Although in many cases we may have a ‘presentation’ of our locally compact quantum group in terms of generators and relations, the generators will usually be unbounded and will have to be understood as operators affiliated with a given operator algebra. This makes the attempts of developing the theory we will discuss in the next two lectures much more difficult.

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